

Nanophysics

(II - MSc PHYSICS - 3rd Semester)

Quantum Confinement

(Unit-I)

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Physics of Quantum Confinement

History: In 1970 Esaki & Tsu proposed fabrication of an artificial structure, consisting of alternating layers of 2 different semiconductors

Layer Thickness

$$\approx 1 \text{ nm} = 10 \text{ \AA} = 10^{-9} \text{ m} \equiv \underline{\text{SUPERLATTICE}}$$

- **PHYSICS**: The main idea was that introduction of an artificial periodicity will “fold” the Brillouin Zones into smaller BZ’s \equiv “**mini-zones**”.

\Rightarrow The idea was that this would **raise the conduction band minima**, which is needed to observe novel physics and can be employed for some quantum and nano device applications.

- **Modern growth techniques** (starting in the 1980's), especially **MBE & MOCVD**, make fabrication of such structures possible!
- For the same reason, it is also possible to fabricate many other kinds of artificial structures on the scale of nm

(**nanometers**) \equiv “**Nanostructures**”

Superlattices = “**2 dimensional**” structures

Quantum Wells = “**2 dimensional**” structures

Quantum Wires = “**1 dimensional**” structures

Quantum Dots = “**0 dimensional**” structures!!

- Clearly, it is not only the electronic properties of materials which can be drastically altered in this way. Also, vibrational properties (phonons). Here, only electronic properties & only an overview!
- For many years, quantum confinement has been a fast growing field in both theory & experiment! It is at the forefront of current research!
- Note that I am not an expert on it!

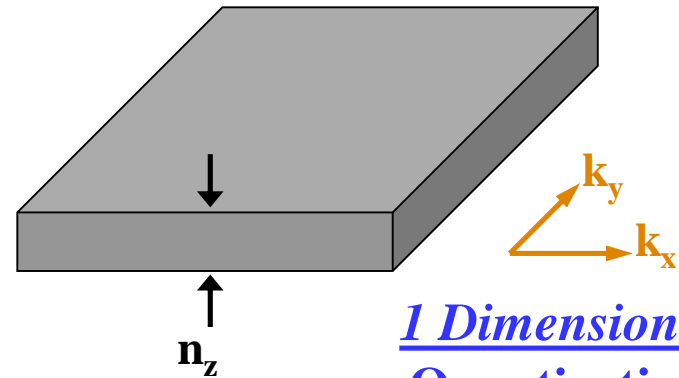
Quantum Confinement in Nanostructures: Overview

Electrons Confined in 1 Direction:

Quantum Wells (thin films):

⇒ **Electrons** can easily move in

2 Dimensions!



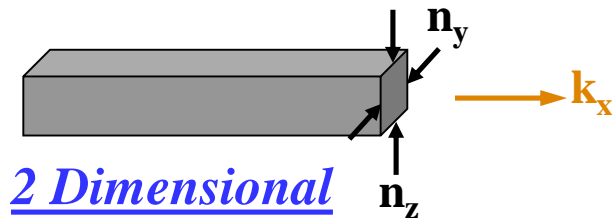
1 Dimensional
Quantization!

Electrons Confined in 2 Directions:

Quantum Wires:

⇒ **Electrons** can easily move in

1 Dimension!



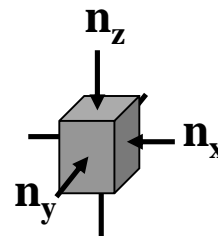
2 Dimensional
Quantization!

Electrons Confined in 3 Directions:

Quantum Dots:

⇒ **Electrons** can easily move in

0 Dimensions!



3 Dimensional
Quantization!

Each further confinement direction changes a continuous k component to a discrete component characterized by a quantum number n .

- **PHYSICS**: Revisit the band structure
 - Consider the 1st Brillouin Zone for the infinite crystal. The maximum wave vectors are of the order

$$\mathbf{k}_m \approx (\pi/a)$$

a = lattice constant. The potential **V** is periodic with period **a**. In the almost free e⁻ approximation, the bands are free e⁻ like except near the Brillouin Zone edge. That is, they are of the form:

$$E \approx (\hbar\mathbf{k})^2/(2m_0)$$

So, the energy at the Brillouin Zone edge has the form:

$$E_m \approx (\hbar\mathbf{k}_m)^2/(2m_0)$$

or

$$E_m \approx (\hbar\pi)^2/(2m_0a^2)$$

PHYSICS

- *SUPERLATTICES* \equiv Alternating layers of material. Periodic, with periodicity \mathbf{L} (layer thickness). Let $\mathbf{k}_z =$ wavevector perpendicular to the layers.
- In a superlattice, the potential \mathbf{V} has a *new periodicity* in the \mathbf{z} direction with periodicity $\mathbf{L} \gg \mathbf{a}$

\Rightarrow In the \mathbf{z} direction, the Brillouin Zone is much smaller than that for an infinite crystal. The maximum wavevectors are of the order:

$$\mathbf{k}_s \approx (\pi/\mathbf{L})$$

\Rightarrow At the BZ edge in the \mathbf{z} direction, the energy has the form:

$$\mathbf{E}_s \approx (\hbar\pi)^2/(2\mathbf{m}_0\mathbf{L}^2) + \mathbf{E}_2(\mathbf{k})$$

$\mathbf{E}_2(\mathbf{k}) =$ the 2 dimensional energy for \mathbf{k} in the \mathbf{x},\mathbf{y} plane.

Note that: $(\hbar\pi)^2/(2\mathbf{m}_0\mathbf{L}^2) \ll (\hbar\pi)^2/(2\mathbf{m}_0\mathbf{a}^2)$

Primary Qualitative Effects of Quantum Confinement

- Consider *electrons confined along 1 direction* (say, **z**) to a layer of width **L**:

Energies

- The *energy bands are quantized* (instead of continuous) in **k_z** & shifted **upward**. So **k_z is quantized**:

$$\mathbf{k}_z = \mathbf{k}_n = [(n\pi)/L], n = 1, 2, 3$$

- So, in the effective mass approximation (**m***), *the bottom of the conduction band is quantized* (like a particle in a **1 d** box) & shifted:

$$\mathbf{E}_n = (n\hbar\pi)^2/(2m^*L^2)$$

- Energies are quantized! Also, the *wavefunctions are 2 dimensional Bloch functions (traveling waves)* for **k** in the **x,y** plane & standing waves in the z direction.

Quantum Confinement Terminology

Quantum Well \equiv QW

= A single layer of material **A** (layer thickness **L**), sandwiched between 2 **macroscopically large layers** of material **B**. Usually, the bandgaps satisfy:

$$E_{gA} < E_{gB}$$

Multiple Quantum Well \equiv MQW

= Alternating layers of materials **A** (thickness **L**) & **B** (thickness **L'**). In this case:

$$L' \gg L$$

So, the e^- & e^+ in one **A** layer are independent of those in other **A** layers.

Superlattice \equiv SL

= Alternating layers of materials **A** & **B** with similar layer thicknesses.

Brief Elementary Quantum Mechanics & Solid State Physics Review

- **Quantum Mechanics of a Free Electron:**

- The **energies are continuous**: $E = (\hbar\mathbf{k})^2/(2m_0)$ (1d, 2d, or 3d)
- The **wavefunctions are traveling waves**:

$$\psi_{\mathbf{k}}(\mathbf{x}) = A e^{i\mathbf{k}\mathbf{x}} \quad (1d)$$

$$\psi_{\mathbf{k}}(\mathbf{r}) = A e^{i\mathbf{k}\cdot\mathbf{r}} \quad (2d \text{ or } 3d)$$

- **Solid State Physics: Quantum Mechanics of an Electron in a Periodic Potential in an infinite crystal :**

- The **energy bands are (approximately) continuous**: $E = E_{\mathbf{n}\mathbf{k}}$
- At the bottom of the conduction band or the top of the valence band, in the effective mass approximation, **the bands can be written:**

$$E_{\mathbf{n}\mathbf{k}} \cong (\hbar\mathbf{k})^2/(2m^*)$$

- The **wavefunctions are Bloch Functions = traveling waves**:

$$\Psi_{\mathbf{n}\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} \mathbf{u}_{\mathbf{n}\mathbf{k}}(\mathbf{r}); \quad \mathbf{u}_{\mathbf{n}\mathbf{k}}(\mathbf{r}) = \mathbf{u}_{\mathbf{n}\mathbf{k}}(\mathbf{r}+\mathbf{R})$$

Some Basic Physics

- Density of states (DoS)**

$$DoS = \frac{dN}{dE} = \frac{dN}{dk} \frac{dk}{dE}$$

in 3D:

$$N(k) = \frac{\text{k space vol}}{\text{vol per state}}$$

$$= \frac{4/3 \pi k^3}{(2\pi)^3 / V}$$

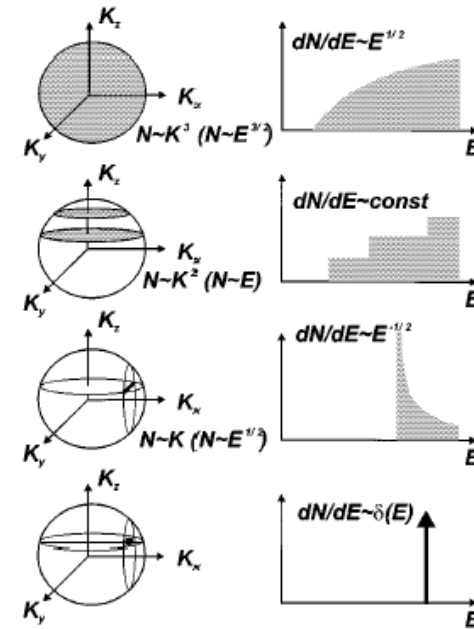


Fig. 1. Density of states for charge carriers in structures with different dimensionalities.

| Structure | Degree of Confinement | $\frac{dN}{dE}$ |
|--------------------|-----------------------|-------------------------------|
| Bulk Material | 0D | \sqrt{E} |
| Quantum Well | 1D | 1 |
| Quantum Wire | 2D | $1/\sqrt{E}$ |
| Quantum Dot | 3D | $\delta(E)$ |

QM Review: The 1D (infinite) Potential Well

(“particle in a box”)

- We want to solve the **Schrödinger Equation for:**

$$\mathbf{x} < \mathbf{0}, \mathbf{V} \rightarrow \infty ; \mathbf{0} < \mathbf{x} < \mathbf{L}, \mathbf{V} = \mathbf{0}; \mathbf{x} > \mathbf{L}, \mathbf{V} \rightarrow \infty$$

$$\Rightarrow -[\hbar^2/(2m_0)](d^2 \psi/dx^2) = E\psi$$

- Boundary Conditions:

$$\psi = 0 \text{ at } \mathbf{x} = \mathbf{0} \ \& \ \mathbf{x} = \mathbf{L} \ (\mathbf{V} \rightarrow \infty \text{ there})$$

- Energies:

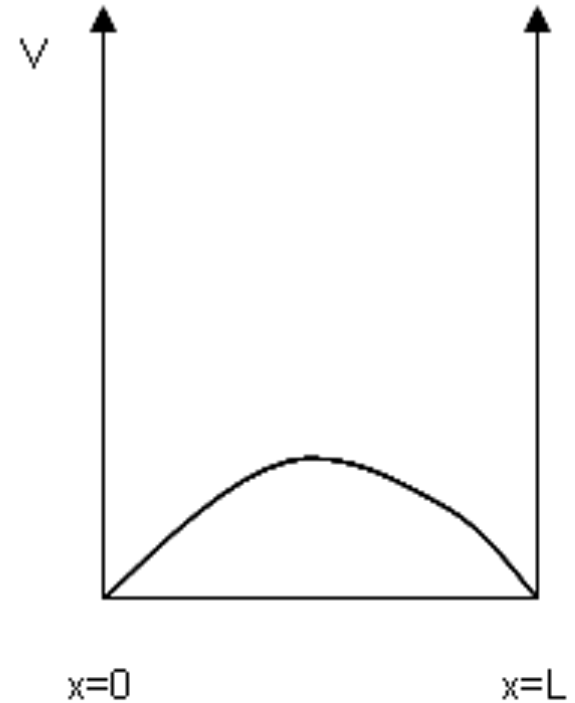
$$\mathbf{E}_n = (\hbar n \pi)^2 / (2m_0 L^2), \quad \mathbf{n} = 1, 2, 3$$

Wavefunctions:

$$\psi_n(\mathbf{x}) = (2/L)^{1/2} \sin(n\pi x/L) \text{ (a standing wave!)}$$

Qualitative Effects of Quantum Confinement:

Energies are quantized & ψ changes from a traveling wave to a **standing wave**.



In 3D_{imensions}...

- For the **3D infinite potential well**:

$$\Psi(x, y, z) \sim \sin\left(\frac{n\pi x}{L_x}\right) \sin\left(\frac{m\pi y}{L_y}\right) \sin\left(\frac{q\pi z}{L_z}\right), \quad n, m, q = \text{integer}$$

$$\text{Energy levels} = \frac{n^2 h^2}{8mL_x^2} + \frac{m^2 h^2}{8mL_y^2} + \frac{q^2 h^2}{8mL_z^2}$$

Real Quantum Structures aren't this simple!!

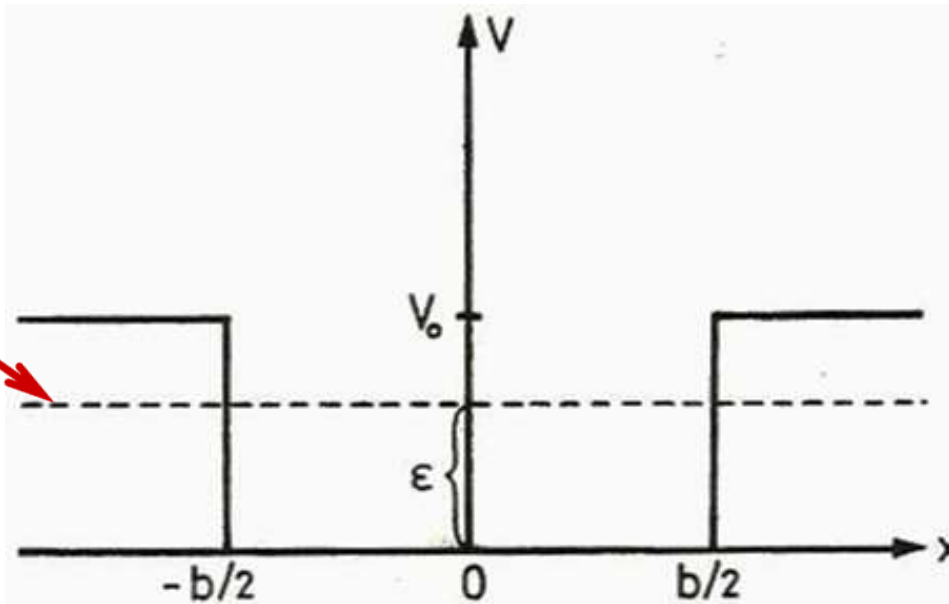
- In **Superlattices & Quantum Wells**, the potential barrier is obviously not infinite!
- In **Quantum Dots**, there is usually ~ **spherical confinement**, not rectangular.
- The simple problem only considers a single electron. But, in **real structures, there are many electrons & also holes!**
- Also, there is often **an effective mass mismatch** at the boundaries. That is *the boundary conditions we've used are too simple!*

QM Review: The 1d (finite) Rectangular Potential Well

In most QM texts!! Analogous to a Quantum Well

- We want to solve the Schrödinger Equation for:

We want bound states: $\epsilon < V_0$



Rectangular potential well

$$[-\{\hbar^2/(2m_0)\}(d^2/dx^2) + V]\psi = \epsilon\psi \quad (\epsilon \equiv E)$$

$$V = 0, \quad -(b/2) < x < (b/2); \quad V = V_0 \text{ otherwise}$$

Solve the Schrödinger Equation:

$$[-\{\hbar^2/(2m_0)\}(d^2/dx^2) + V]\psi = \epsilon\psi$$

$$(\epsilon \equiv E) \quad V = 0, \quad -(b/2) < x < (b/2)$$

$$V = V_0 \text{ otherwise}$$

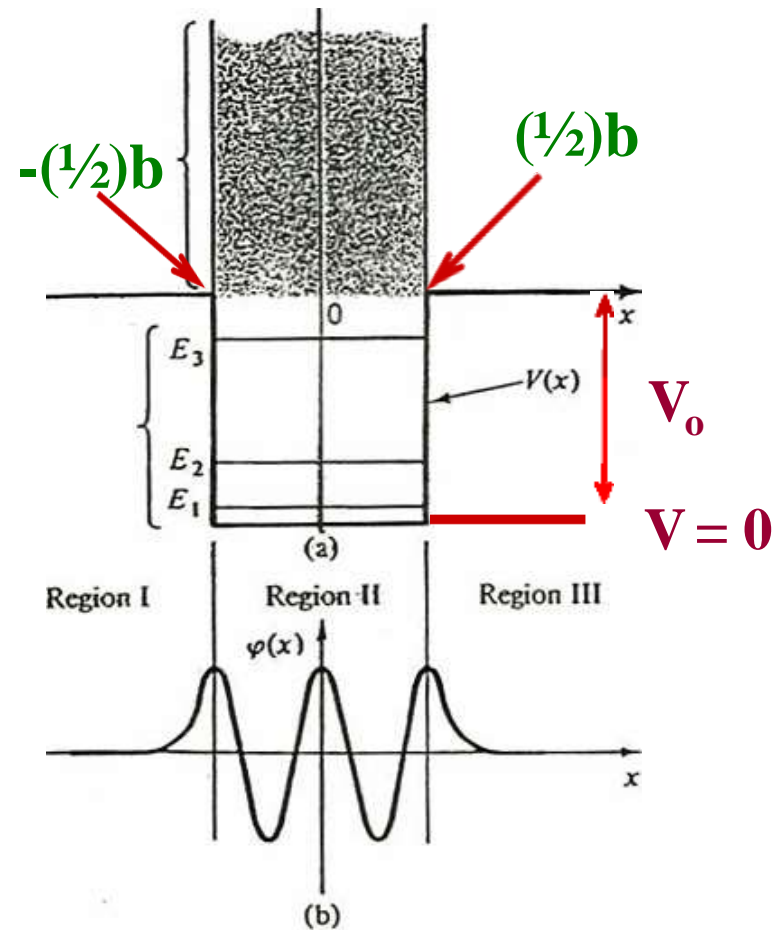
Bound states are in Region II

Region II:

$\psi(x)$ is oscillatory

Regions I & III:

$\psi(x)$ is decaying



Finite rectangular potential well. (a) The potential function $V(x)$ and energy spectrum. (b) Typical structure of a bound eigenstate. Function oscillates in region II where kinetic energy is positive and decays in regions I and III, where kinetic energy is negative.

The 1d (finite) rectangular potential well

A brief math summary!

Define: $\alpha^2 \equiv (2m_0\varepsilon)/(\hbar^2)$; $\beta^2 \equiv [2m_0(\varepsilon - V_0)]/(\hbar^2)$

The Schrödinger Equation becomes:

$$(d^2/dx^2) \psi + \alpha^2 \psi = 0, \quad -(1/2)b < x < (1/2)b$$

$$(d^2/dx^2) \psi - \beta^2 \psi = 0, \quad \text{otherwise.}$$

⇒ Solutions:

$$\psi = C \exp(i\alpha x) + D \exp(-i\alpha x), \quad -(1/2)b < x < (1/2)b$$

$$\psi = A \exp(\beta x), \quad x < -(1/2)b$$

$$\psi = A \exp(-\beta x), \quad x > (1/2)b$$

Boundary Conditions:

⇒ ψ & $d\psi/dx$ are continuous SO:

- Algebra (2 pages!) leads to:

$$(\epsilon/V_0) = (\hbar^2\alpha^2)/(2m_0V_0)$$

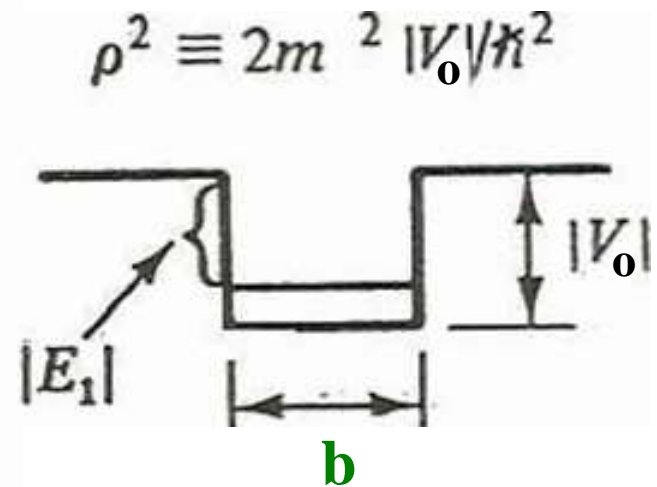
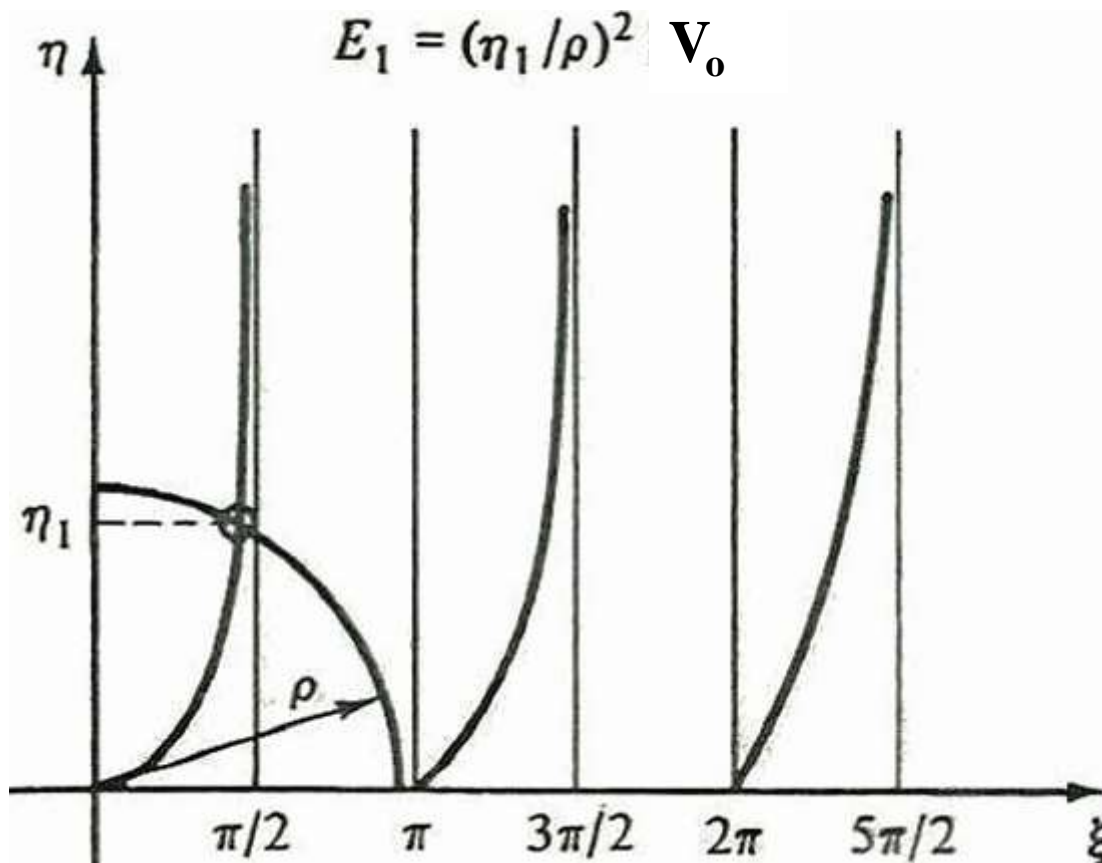
ϵ , α , β are related to each other by transcendental equations.

For example:

$$\tan(\alpha b) = (2\alpha\beta)/(\alpha^2 - \beta^2)$$

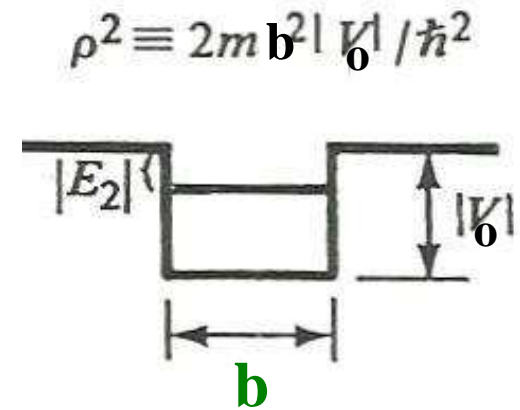
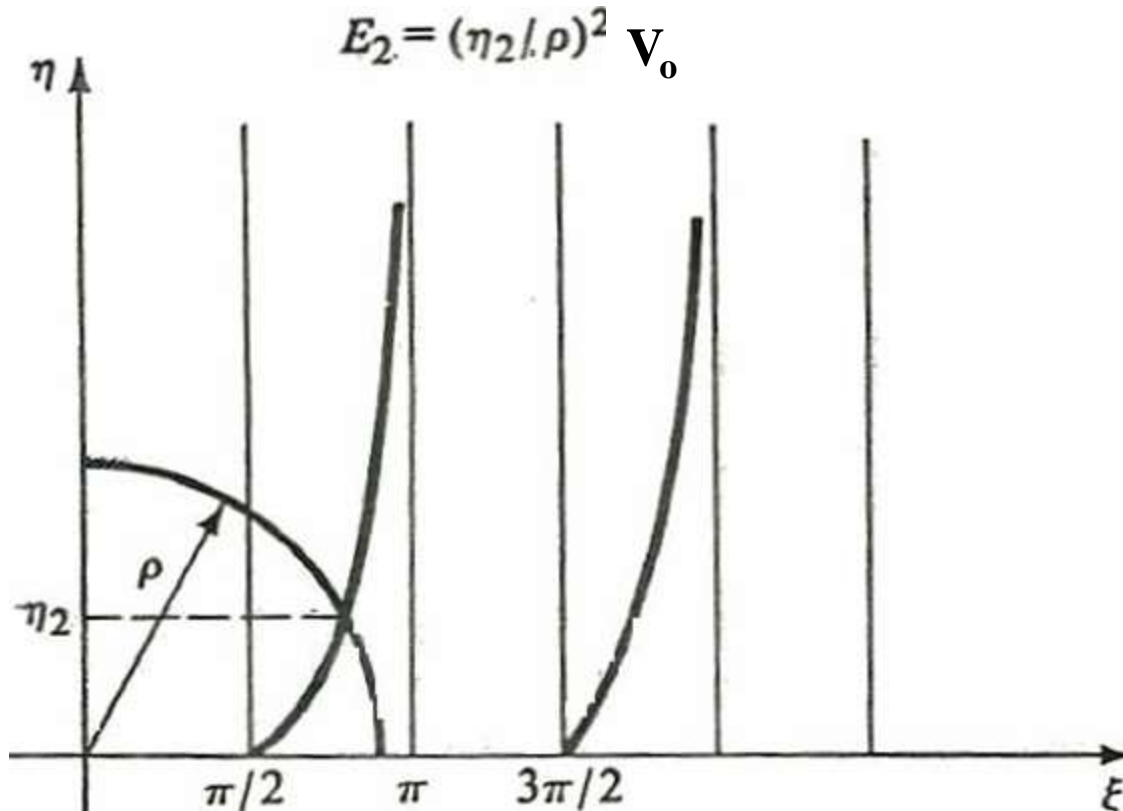
- Solve graphically or numerically.
- **Get:** *Discrete Energy Levels* in the well
(a finite number of finite well levels!)

- Even eigenfunction solutions (a finite number):
Circle, $\xi^2 + \eta^2 = \rho^2$, **crosses** $\eta = \xi \tan(\xi)$



- Odd eigenfunction solutions:

Circle, $\xi^2 + \eta^2 = \rho^2$, **crosses** $\eta = -\xi \cot(\xi)$



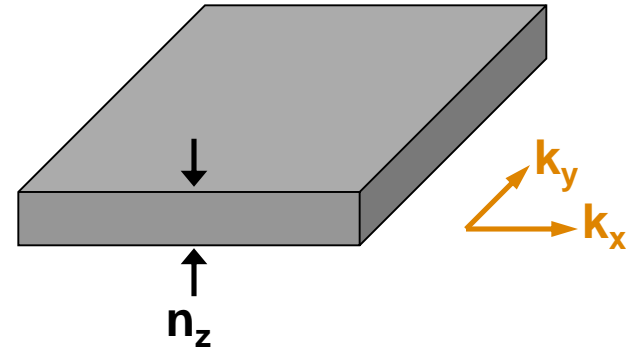
$$|E_2| < |E_1|$$

Quantum Confinement in Nanostructures

Confined in:

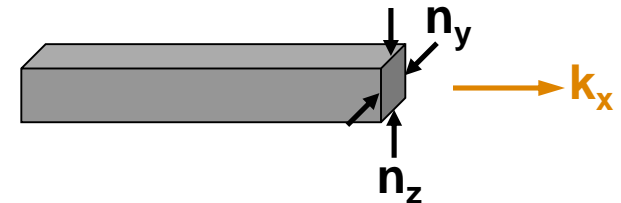
1 Direction: Quantum well (thin film)

Two-dimensional electrons



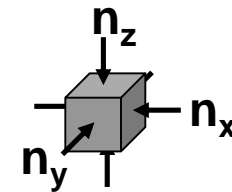
2 Directions: Quantum wire

One-dimensional electrons



3 Directions: Quantum dot

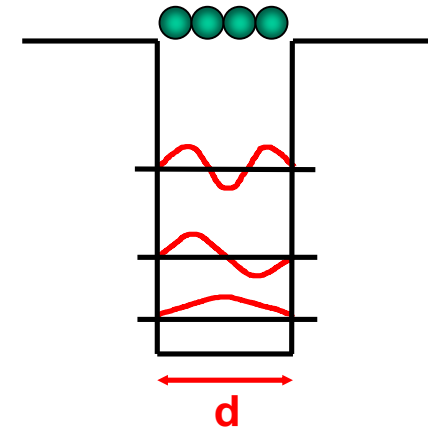
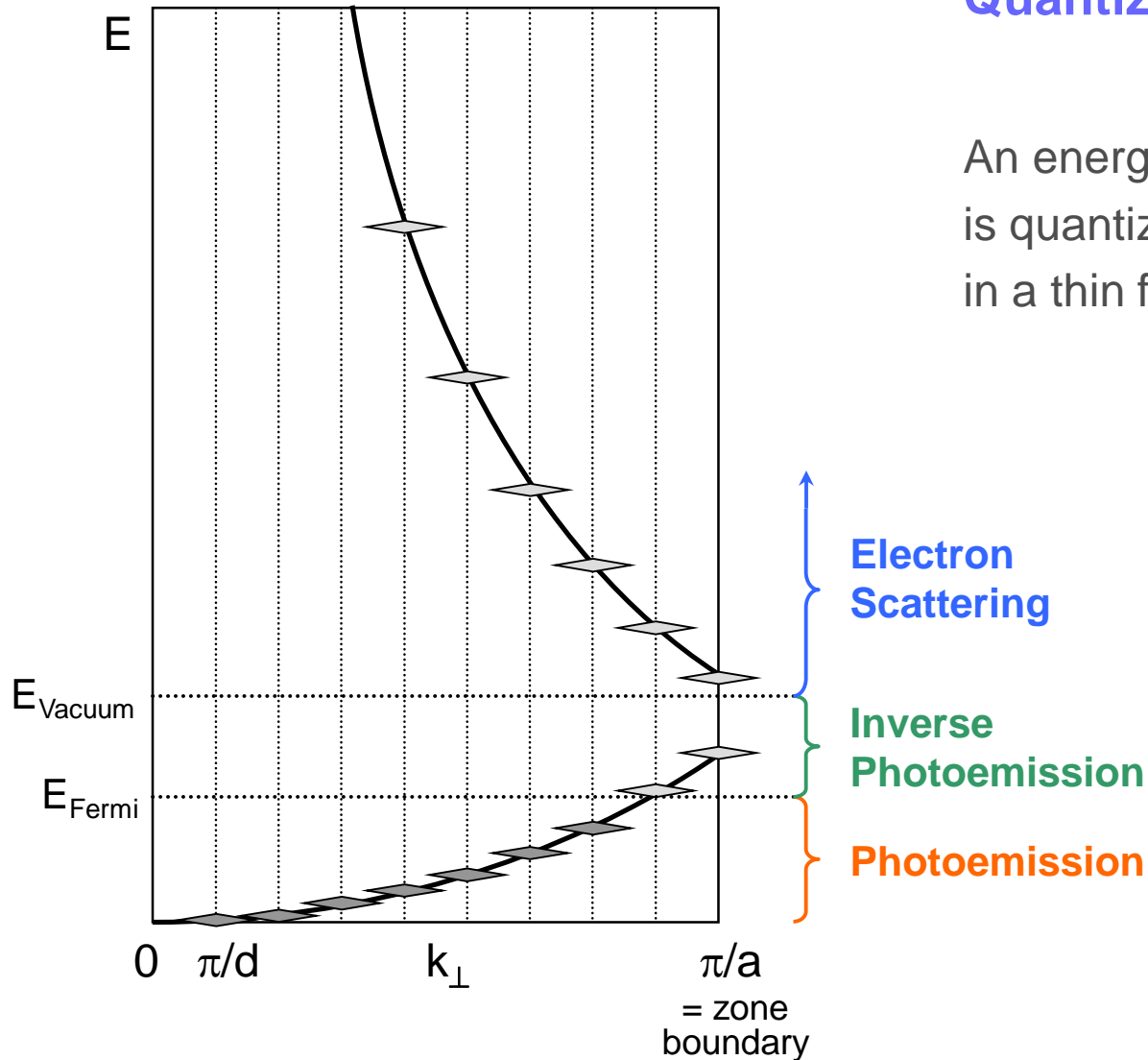
Zero-dimensional electrons



Each confinement direction converts a continuous k in a discrete quantum number n .

Quantization in a Thin Crystal

An energy band with continuous k is quantized into N discrete points k_n in a thin film with N atomic layers.



N atomic layers with the spacing $a = d/n$

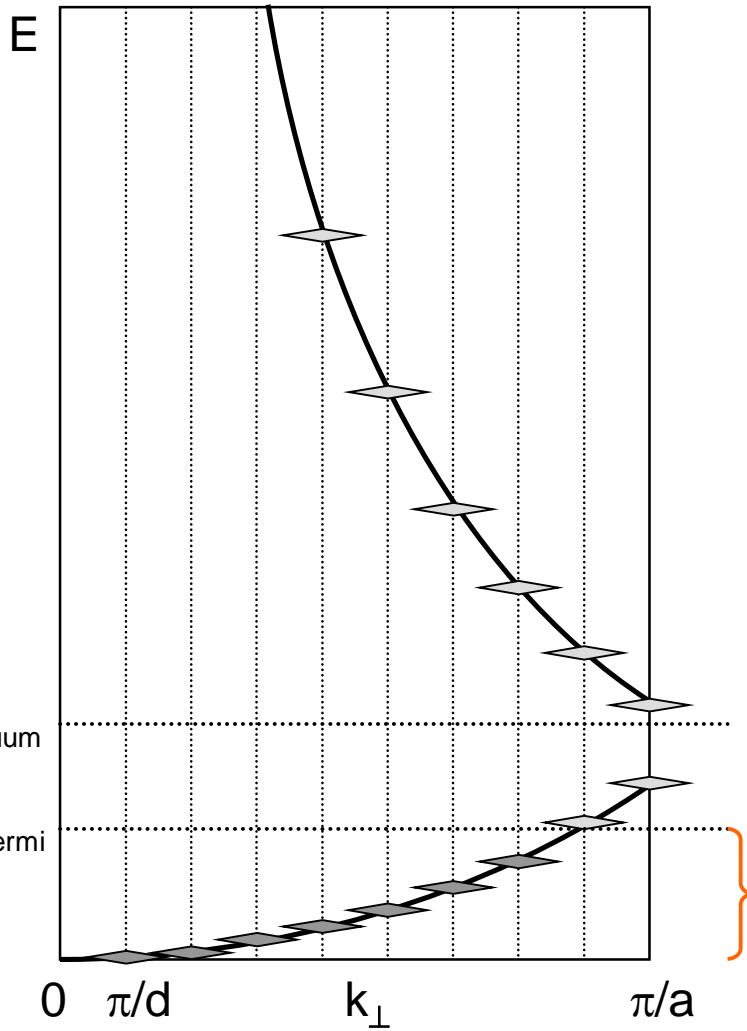
N quantized states with $k_n \approx n \cdot \pi/d$ ($n = 1, \dots, N$)

$$\lambda_n = 2d/n$$

$$k_n = 2\pi/\lambda_n = n \cdot \pi/d$$

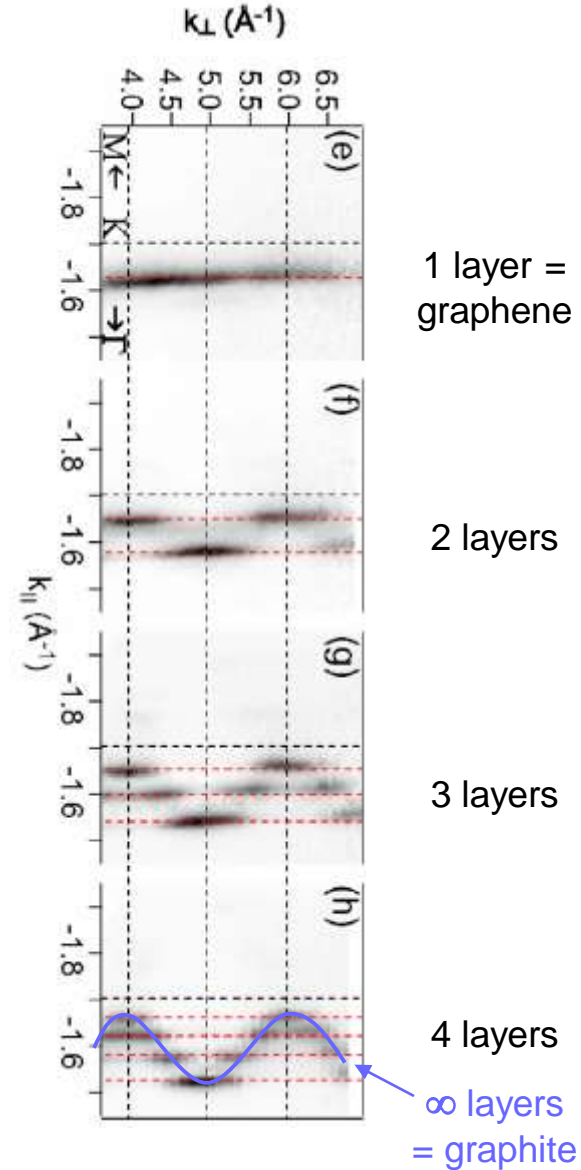
Quantization in Thin Graphite Films

Lect. 7b,
Slide 11



Photoemission

N atomic layers with spacing $a = d/n$:
 \Rightarrow **N** quantized states with $k_n \approx N \cdot \pi/d$



1 layer =
graphene

2 layers

3 layers

4 layers

∞ layers
= graphite

The Important Electrons in a Metal

$$\text{Energy} \approx E_{\text{Fermi}}$$

$$\text{Energy Spread} \approx 3.5 k_{\text{B}}T$$

Transport (conductivity, magnetoresistance, screening length, ...)

Width of the Fermi function:

$$\text{FWHM} \approx 3.5 k_{\text{B}}T$$

Phase transitions (superconductivity, magnetism, ...)

Superconducting gap:

$$E_{\text{g}} \approx 3.5 k_{\text{B}}T_{\text{c}} \quad (T_{\text{c}} = \text{critical temperature})$$